Controlling Intermolecular Associations with Molecular Superstructure: Polar Discotic Linear Chain Phases

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The design of new molecule-based materials which exhibit specific properties requires a careful balance of both intermolecular interactions and molecular superstructure. The generation of polar molecule-based materials is a particularly important objective since it is a necessary prerequisite for ferroelectricity¹ and second-order nonlinear optical susceptibility,² properties of immense technological importance. However, it is generally difficult to generate thermodynamically stable polar assemblies due to the tendency of materials to avoid internal electric fields by adopting antiferroelectric structures. To overcome this obstacle and find an optimal balance of interactions which will generate an achiral ferroelectric liquid crystalline state,^{3,4} we have been interested in hexagonal columnar superstructures with axial polarity since the triangular symmetry of the lattice cannot accommodate bulk antiferroelectric order.5 We report herein a number of new hexagonal and rectangular columnar liquid crystalline phases (mesophases) with axial polarity resulting from head-to-tail alignment of vanadyl groups in linear chain structures, $(-V=O-V=O-)_n$. Further, we demonstrate the use of design principles whereby intermolecular associations may be controlled in a predictable fashion by designing correlated molecular superstructures.

A series of liquid crystalline vanadyl Schiff-base complexes 1, 2, and 3 (Scheme I) have been investigated which exhibit different degrees of intermolecular associations via dative bonding between vanadyl groups. The relative strength of dative bonding and subsequent strength of the linear chain structure $(-V=O-V=O-)_n$ may be determined qualitatively by analyzing the frequency of the V=O stretching mode, wherein stronger linear chains are observed at lower frequency.⁶ Consistent with our previous investigations on smectic liquid crystals,⁶ all of the materials are monomeric in their isotropic phases (V=O 992-994 cm⁻¹), 1b, 2b, and 3b tend to exhibit the strongest linear chains (V=O 854-868 cm⁻¹), 1c, 2c, and 3c display somewhat weaker linear chains (V=O 871-911 cm⁻¹), and 1a, 2a, and 3a are monomeric (V=O 992 cm⁻¹). The electronic transitions of 1, 2, and 3 are insensitive to the nature of the mesophase; however, in monomeric form these complexes are green with broad metal-centered transitions at 401-404 nm, and in linear chain structures they are orange with transitions at 347-355 nm.

The complexes in Scheme I were also chosen to produce columnar mesophases.⁷ A hexagonal organization of columns

(6) Serrette, A.; Carroll, P. J.; Swager, T. M. J. Am. Chem. Soc. 1992, 114, 1887.

 (D_h) implies that the molecules are perpendicular to and project a time-averaged circular shape along the column axis. In rectangular (D_r) or oblique (D_{ob}) columnar mesophases, the molecules project an elliptical shape and are generally tilted with respect to the column's normal. Complexes with general structures 1 and 2 do not have the disc-like shape generally required for a columnar liquid crystal, and hence columnar phases of these molecules exhibit the correlated structures shown in Scheme I.⁸ The complexes in series 3, with twice the number of side chains, approximate a disc-like shape in monomeric form, and hence columnar phases of these compounds have less correlation.

With the exception of 2c, all of the complexes display enantiotropic mesophases as shown in Table I. The phase assignments were made through a combination of X-ray diffraction (XRD),⁹ optical textures, and miscibility studies.⁷ Phases labeled as disordered ($D_{rd}, D_{hd}, D_{hdl},$ etc.) exhibited XRD patterns with broad amorphous halos at wide angles, indicating liquidlike order of the mesogens along the column axis, and phases labeled as ordered (D_{ro}, D_{ho}) displayed wide-angle peaks consistent with a more regular period along the column axis.

The monomeric complex 1a exhibits a single D_{hd} phase, whereas the linear chain derivatives 1b and 1c display a number of additional unusual hexagonal phases (D_{hd1}, D_{hd2}, D_{hd3}, and D_{ho1}). 1a displays a D_{hd} XRD pattern (a = 44.7 Å) consisting of only a strong (100) peak indicating a fairly disordered phase. However, the mesophases of 1b and 1c display a number of stronger and higher order diffraction peaks.9 These polymeric phases still exhibit liquid-like order within the columns, and the higher order XRD peaks index to a 2D hexagonal lattice, indicating a more regular organization of the columns.¹⁰ The antiparallel correlation in 1c's superstructure positions the geminal methyls in the clefs presented by nearest-neighbor complexes. This lock-and-key organization restricts the mobility of 1c's mesophases and results in a much larger isotropic transition enthalpy than is observed for 1b. The higher correlation in 1c's phases produces larger intercolumnar spacings which range from 46.8 Å in the D_{hd1} to 49.0 Å in the D_{hd3} , whereas those of 1b range from 44.6 Å in the D_{hd1} to 45.8 Å in the D_{ho1} . The V=O stretching bands of both 1b and 1c shift with cooling by 4 and 17 cm⁻¹, respectively, to lower energy, indicating a strengthening of the linear chain. The larger range of the V=O stretching band of 1c is also a consequence of the steric demands of geminal methyls since torsional motions between mesogens generate steric repulsions.

2a exhibits a single D_{hd} mesophase (a = 46.5 Å), whereas the linear chain analog **2b** displays a D_{hd} phase (a = 43.2 Å) at high temperature, followed by rectangular disordered and crystal (D_{rd} and K) phases at lower tempertures. The similarities of these lattice constants with those of 1a and 1b are consistent with the correlated superstructures shown in Scheme I.8 2c does not display a stable mesophase since the superstructure shown in Scheme I directs the geminal methyls toward the aromatic rings and thereby introduces steric repulsions which reduce dipolar and dative interactions. In the case of 2b, the D_{hd} -to- D_{rd} transition is accompanied by an abrupt shift of the V=O band from 855 to 861 cm⁻¹ and a large decrease in birefringence. This weakening of the linear chain and the change in birefringence in 2b's D_{rd}-(P21/a) phase indicates a substantial tilting of the mesogenic cores which reduces the interaction between the apical oxos and neighboring vanadium centers.

[†] Office of Naval Research Young Investigator, 1992–1995; NSF-Young Investigator, 1992–1997.

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⁽³⁾ Ferroelectricity has been studied extensively in chiral liquid crystals (ref 1). However, the achiral liquid crystals there has been only one report of ferroelectricity (ref 4).

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A. M.; Malthete, J. Mol. Cryst. Liq. Cryst. 1984, 106, 121. (b) Billard, J.
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Chemical Physics; Springer: Berlin, 1980, p 383. (c) Chandrasekhar, S.;
Ranganath, G. S. Rep. Prog. Phys. 1990, 53, 57 and references therein.
(8) For a previous demonstration of this concept, see: Lai, C. K.; Serrette,

⁽a) For a previous demonstration of this concept, see. Lat, C. K., Serrette, A.; Swager, T. M. J. Am. Chem. Soc. 1992, 114, 7948-7949.

⁽⁹⁾ XRD data are given in the supplementary material.

⁽¹⁰⁾ Similar XRD patterns have been reported for other metallomesogens: Ohta, K.; Hasebe, H.; Moriya, M.; Fujimoto, T.; Yamamoto, I. Mol. Cryst. Liq. Cryst. 1991, 208, 43.



Table I. Phase Assignments, Transition Temperatures (°C, above arrow), and Transition Enthalpies (kcal mol, below arrows, in parentheses) for All Compounds Described (The range of the V=O stretching band frequency (cm^{-1}) is given below the phase designation.)

1a	$D_{hd} = \frac{150}{(5.3)}$	$\frac{1}{1} \rightarrow I$			
16	$D_{ho1} = \frac{119}{(4.5)}$ 861-862	$D_{hd3} \rightarrow D_{hd3} \rightarrow \frac{13}{(2)}$ 862-863	$ \begin{array}{c} \underline{56.1} \rightarrow D_{hd2} \rightarrow \underline{16} \\ \underline{5.5} \qquad 863 \end{array} $	$D_{hd1} = D_{hd1} = \frac{17}{(0)}$ 863-865	(<u>5.0</u> ►I (.76) ►I 994
1 c	K -27. 871	$\frac{7}{9} \rightarrow D_{hd3} \rightarrow \frac{66}{(3.4)}$ 871-873	$\frac{1.6}{4} \rightarrow D_{hd2} \rightarrow \frac{91}{(4)}$ 873-883	$\frac{.1}{(11)} \rightarrow D_{hd1} \rightarrow \frac{146}{(11)}$ 883-888	8 (6) ► I 992
2a	D _{hd} = 155.8 (0.71) 992	³ → I 992			
2b	K - 62.1 (20.7) 861*	$D_{rd1} \rightarrow \frac{107.}{(2.2)}$ 858-861*	$D_{hd} = \frac{169.7}{(1.0)}$ 855-861	→ I 994	
2c	K 78.3 (18.3) 865	M ^b 114.8 (8.1) 865-875	- I 992		
3a	D _{ro} <u>119</u> (4.4 992	$D_{rd} = \frac{156.1}{(2.5)}$ 992	→ I 992		
3b	D _{ho} ≼ 36. 864-868	$\frac{4}{(2.6)}$ D _{hd} $\frac{97.5}{(2.6)}$ 868-854	$D_{rd} = \frac{151.4}{(2.5)}$ 854-856	→ I 994	
3c	D _{hd} <u>55.3</u> 911 ^a	$D_{rd} = \frac{107.8}{(8.7)}$ 904-889	► I 992		

^a The stretching band changes very abruptly at the phase transition. ^b Transient mesophase observed only on first heating.

The greater side-chain density in 3a, 3b, and 3c destabilizes the hexagonal structure, and all three compounds exhibit D_{rd} phases at higher temperatures. 3c has a weakened linear chain and a lower clearing point than 3a and 3b, which is consistent with steric interactions between the randomly oriented geminal dimethyl groups and neighboring mesogens. Interestingly, both 3b and 3c exhibit an unusual inverted phase behavior¹¹ with the more ordered D_{rd} phase at higher temperatures, followed by a less ordered D_{hd} phase at lower temperatures. The intercolumnar spacing of 3b's D_{hd} phase is 51.0 Å, which is very similar to those of 1c, suggesting that the steric locking in 1c's mesophases produces similar spatial requirements. In contrast to what was observed for 2b, the linear chains of 3b and 3c are stronger in the D_{rd} phases than in the D_{hd} phases. The origin of this unexpected result is unclear; however, these D_{rd} phases have C2/m symmetry and exhibit similar birefringence to that of the D_{hd} phases, suggesting that the D_{rd} phase is only slightly tilted.

In summary, we have demonstrated that liquid crystalline superstructures can be designed which selectively control the degree of intermolecular associations. In general, intermolecular dative interactions and the resulting polar order produce a richer liquid crystalline behavior.¹² Linear chain D_{hd} phases were observed which gave XRD profiles typical of monomeric liquid crystals. The linear chains in these phases are partly melted and/or weak and thereby allow some transverse mobility of the mesogens and less ordered columns. Linear chain structures can also produce highly organized columnar structures in which the mesogens are held in tight registry. We find that depending on their nature, rectangular phases can weaken or strengthen the dative bonding between vanadyl groups relative to a hexagonal organization.

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Supplementary Material Available: Plots of V=O stretching frequency as a function of temperature, XRD data, and infrared data for V= 18 O-labeled compounds (9 pages). Ordering information is given on any current masthead page.

⁽¹¹⁾ For other examples of inverted behavior in discotics, see: Destrade, C.; Gasparoux, H.; Babeau, A.; Nguyen, H. T. Mol. Cryst. Liq. Cryst. 1981, 67, 37.

⁽¹²⁾ Monomeric discotics with dipoles normal to the mesogens plane have been previously found to exhibit antiparallel dipolar organization. Piechocki, C.; Boulou, J.-C.; Simon, J. Mol. Cryst. Liq. Cryst. 1987, 149, 115.